

Integral Equations and Fast Algorithms

Lecture 16: Helmholtz

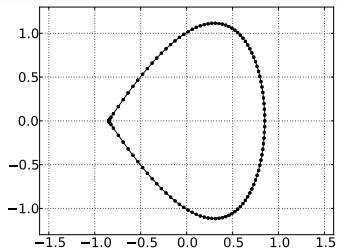
CS 598AK · October 17, 2013

Outline

Harmonic Potential Theory

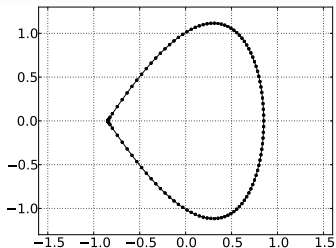
Helmholtz Potential Theory

Domains with corners



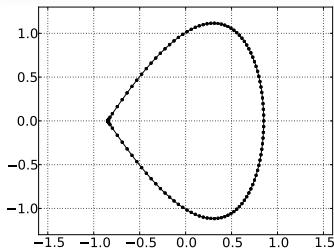
What's the problem?

Domains with corners



What's the problem? (*Hint: Jump condition for constant density*)

Domains with corners

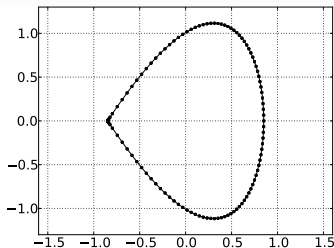


What's the problem? (*Hint: Jump condition for constant density*)

At corner x_0 : (2D)

$$\lim_{x \rightarrow x_0 \pm} = PV \int_{\partial\Omega} \hat{n} \cdot \nabla_y G(x, y) \varphi(y) ds_y \pm \frac{1}{2} \frac{\langle \text{opening angle on } \pm \text{ side} \rangle}{\pi} \varphi$$

Domains with corners



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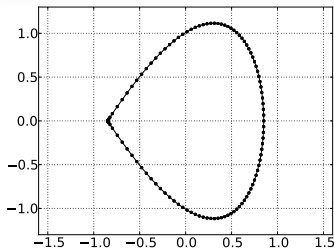
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→ non-continuous behavior of potential on Γ at x_0

What space have we been living in?

Domains with corners



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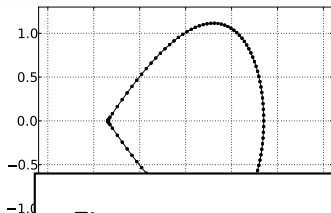
→ non-continuous behavior

What space have we been

Fixes:

- $I +$ Bounded (Neumann) + Compact (Fredholm)
- Use L^2 theory (point behavior “invisible”)

Domains with corners



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Numerically: Needs consideration, but ultimately straightforward to fix. (once you know how)

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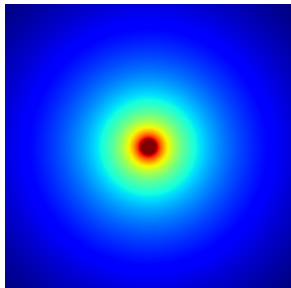
Harmonic Potential Theory

Helmholtz Potential Theory

Helmholtz

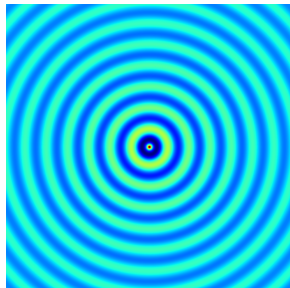
Laplace Equation

$$-\Delta u = \delta$$



Helmholtz Equation

$$\Delta u + k^2 u = \delta$$

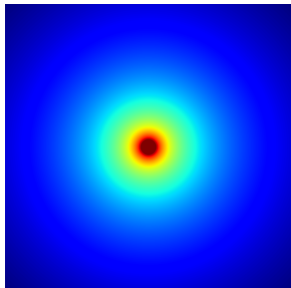


(where $k = \omega/c$)

Helmholtz

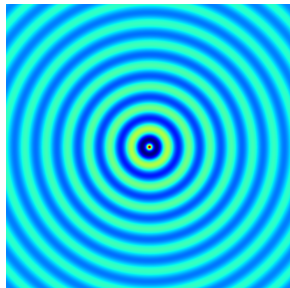
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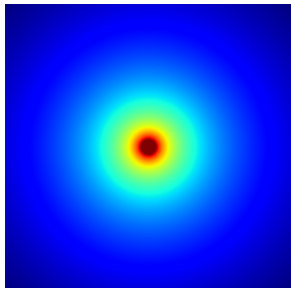


What happens for increasing k ?

Helmholtz

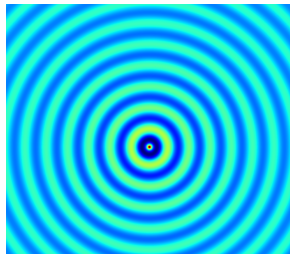
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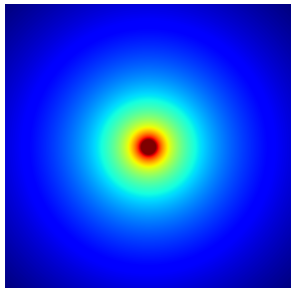
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...and complex k ?

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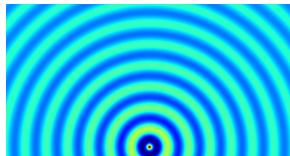
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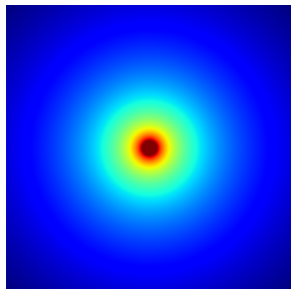
...and complex k ?

What is c and where does it come from?

Helmholtz

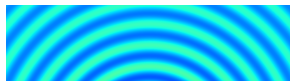
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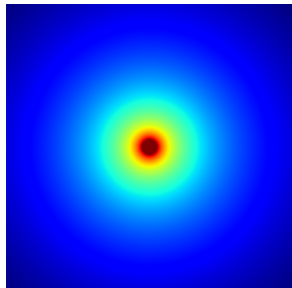
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$$\partial_t^2 u = c^2 \Delta u$$

Helmholtz

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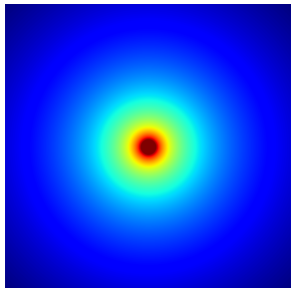
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How to recover the time-dependent solution?

Helmholtz

Laplace Equation

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Helmholtz Equation

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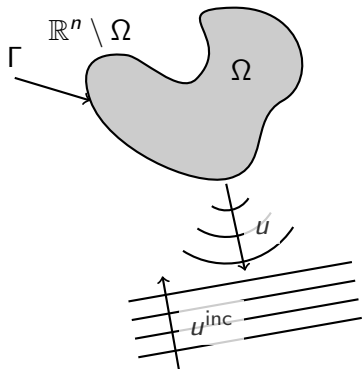
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How to recover the time-dependent solution?

Typical problems being solved this way?

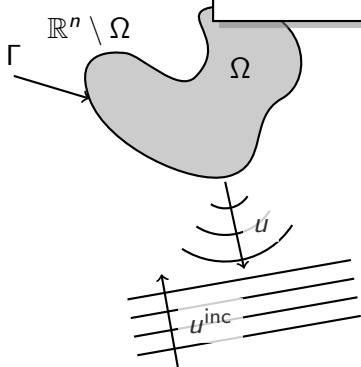
Scattering problem



$$u^{tot} = u + u^{inc}$$

Scattering problem

Where do the two couple?



$$u^{\text{tot}} = u + u^{\text{inc}}$$

Back to physics

Physical quantities:

- Velocity potential: $U(x, t) = u(x)e^{-i\omega t}$
- Velocity: $v = (1/\rho_0)\nabla U$
- Pressure: $p = -\partial_t U = i\omega u e^{-i\omega t}$
 - Equation of state: $p = f(\rho)$

(fix phase by e.g. taking real part)

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What's ρ_0 ?

What happens to a pressure BC as $\omega \rightarrow 0$?

Fundamental solutions

$$G_k(x) = \begin{cases} \frac{i}{4} H_0^{(1)}(k|x|) & 2\text{D} \\ \frac{1}{4\pi} \frac{e^{ik|x|}}{|x|} & 3\text{D} \end{cases}$$

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Hankel function?

Separation of variables in polar coordinates \rightarrow Bessel ODE:

$$x^2 \frac{d^2 y}{dx^2} + x \frac{dy}{dx} + (x^2 - \alpha^2)y = 0$$

$$H_\alpha^{(1)} = J_\alpha + iY_\alpha,$$

$$H_\alpha^{(2)} = J_\alpha - iY_\alpha.$$

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Asymptotic behavior for small k ?

$G_k(x)/G_0(x) \rightarrow ?$ ($x \rightarrow 0$)

What happens for $-k$?

Helmholtz Boundary Conditions

- **Sound-soft:** Pressure remains constant

- Scatterer “gives”

$$u = f \rightarrow \text{Dirichlet}$$

- **Sound-hard:** Pressure same on both sides of interface

- Scatterer “does not give”

$$\hat{n} \cdot \nabla u = 0 \rightarrow \text{Neumann}$$

- **Impedance:** Some pressure translates into motion

- Scatterer “resists”

$$\hat{n} \cdot \nabla u + ik\lambda u = 0 \rightarrow \text{Robin } (\lambda > 0)$$

- **Sommerfeld** radiation condition: allow only outgoing waves

$$(r \rightarrow \infty)$$

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$$r^{\frac{n-1}{2}} \left(\frac{\partial}{\partial r} - ik \right) u(x) \rightarrow 0 \quad (r \rightarrow \infty)$$

Helmholtz Boundary Conditions

- **Sound-soft:** Pressure = 0. Many interesting BCs → many IEs! :)
 - Scatterer “gives” $u = f \rightarrow$ Dirichlet
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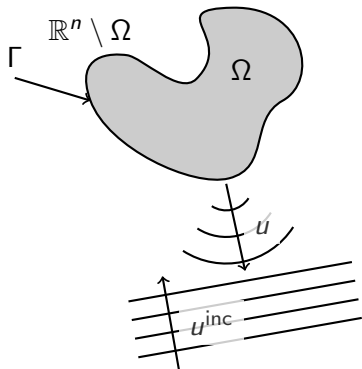
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Many interesting BCs \rightarrow many IEs! :)

Transmission between media: What's continuous?

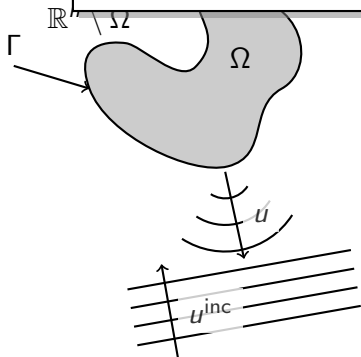
Scattering problem



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Scattering problem

What's allowed as an incoming wave?



$$u^{tot} = u + u^{inc}$$

Unchanged from Laplace

Theorem (Green's Formula [Colton/Kress IAEST Thm 2.1])

If $\Delta u + k^2 u = 0$, then

$$(S(\hat{n} \cdot \nabla u) - Du)(x) = \begin{cases} u(x) & x \in D \\ \frac{u(x)}{2} & x \in \partial D \\ 0 & x \notin D \end{cases}$$

$$\begin{aligned} \lim_{x \rightarrow x_0 \pm} (S'u) &= \left(S' \mp \frac{1}{2} I \right) (u)(x_0) &\Rightarrow [Su] &= 0 \\ \lim_{x \rightarrow x_0 \pm} (Du) &= \left(D \pm \frac{1}{2} I \right) (u)(x_0) &\Rightarrow [S'u] &= -u \\ & &\Rightarrow [Du] &= u \\ & &\Rightarrow [D'u] &= 0 \end{aligned}$$

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Why is singular behavior (esp. jump conditions) unchanged?

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$$e^{ikr} = 1 + O(r)$$

Green's thm reminder:

$$\int_{\Omega} u \Delta v - v \Delta u = \int_{\partial \Omega} u(\hat{n} \cdot \nabla v) - v(\hat{n} \cdot \nabla u) ds$$

$\lim_{x \rightarrow x_0}$

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Resonances

– Δ on a bounded (interior) domain with homogeneous Dirichlet/Neumann BCs has countably many real, positive eigenvalues.

Problem?

Resonances

– Δ on a bounded (interior) domain with homogeneous Dirichlet/Neumann BCs has countably many real, positive eigenvalues.

Problem?

$$-\Delta u = \lambda u$$

$$\Delta u + k^2 u = 0$$

Questions?

?