

## Integral Operators are Compact

**Theorem 15.** (Continuous kernel  $\Rightarrow$  compact [Kress LIE Thm. 2.21])  $G \subset \mathbb{R}^n$  compact,  $K \in C(G^2)$ . Then

$$(A\phi)(x) := \int_G K(x, y) \phi(y) dy.$$

is compact on  $C(G)$ .

\* Use A-A. (a statement about compact [sets](#))

What is there to show?

- Pick  $U \subset C(G)$ .  $A(U)$  bounded?
- $A(U)$  equicontinuous?

## Weakly singular

$G \subset \mathbb{R}^n$  compact

### Definition 16. (Weakly singular kernel)

- $K$  defined, continuous everywhere except at  $x = y$
- There exist  $C > 0$ ,  $\alpha \in (0, n]$  such that

$$|K(x, y)| \leq C |x - y|^{\alpha - n} \quad (x \neq y)$$

### Theorem 17. (Weakly singular kernel $\Rightarrow$ compact [Kress LIE Thm. 2.22])

$K$  weakly singular. Then

$$(A\phi)(x) := \int_G K(x, y) \phi(y) dy.$$

is compact on  $C(G)$ .

- Outline the proof.

## Weakly singular (on surfaces)

- $\Omega \subset \mathbb{R}^n$  bounded, open,  $C^1$

**Definition 18. (Weakly singular kernel (on a surface))**

- $K$  defined, continuous everywhere except at  $x = y$
- There exist  $C > 0$ ,  $\alpha \in (0, n - 1]$  such that

$$|K(x, y)| \leq C |x - y|^{\alpha - n + 1} \quad (x, y \in \partial \Omega, x \neq y)$$

**Theorem 19. (Weakly singular kernel  $\Rightarrow$  compact [Kress LIE Thm. 2.23])**  $K$  weakly singular on  $\partial \Omega$ . Then

$$(A\phi)(x) := \int_G K(x, y) \phi(y) dy.$$

is compact on  $C(G)$ .

## 6.4 Riesz and Fredholm

## Riesz Theory (I)

- Still trying to solve

$$L\phi := (I - A)\phi = \phi - A\phi = f$$

with  $A$  compact.

**Theorem 20. (First Riesz Theorem [Kress, Thm. 3.1])**  $N(L)$  is finite-dimensional.



Questions:

- What is  $N(L)$  again?
- Why is this good news?
- Show it.

## Riesz Theory (Part II)

**Theorem 21. (Riesz theory [Kruskal, Thm. 3.4])** *A compact. Then:*

- $(I - A)$   $\overset{\text{(I)}}{\text{injective}} \Leftrightarrow (I - A) \overset{\text{(II)}}{\text{surjective}}$ 
  - *It's either bijective or neither s nor i.*
- *If  $(I - A)$  is bijective,  $(I - A)^{-1}$  is bounded.*

◦ Rephrase for solvability

$$\text{(I)} \quad \varphi - \int \kappa \varphi = 0 \Rightarrow \varphi = 0$$

◦ Main impact?

$$\text{(II)} \quad \forall f \in C: \varphi - \int \kappa \varphi = f$$

◦ Key shortcoming?

## Hilbert spaces

Hilbert space: Banach space with a norm coming from an inner product:

$$(\alpha x + \beta y, z) = ?$$

$$(x, \alpha y + \beta z) = ?$$

$$(x, x) \stackrel{0 \in \mathbb{R}}{=} x = 0$$

$$(y, x) = ?$$

$$(f, g) = \int f(x) \overline{g(x)} dx$$

$$\left(\vec{x}, \vec{y}\right) = \sum_i x_i y_i$$

$$\|f\|_2 = \sqrt{\langle f, f \rangle}$$

$\mathbb{C}^2$

$$\|f\|_\infty = \sup_{x \in G} |f(x)|$$

- Is  $C^0(G)$  a Hilbert space?
- Name a Hilbert space of functions.  $\longrightarrow$
- Is  $C^0(G)$  "equivalent" to  $L^2(G)$ ?
- Why do compactness results transfer over nonetheless? Hint: What is

$$|(f, g)| \leq \|f\|_2 \|g\|_2$$

## Adjoint Operators

**Definition 22. (Adjoint operator)**  $A^*$  called *adjoint to A* if

$$(Ax, y) = (x, A^*y)$$

for all  $x, y$ .

Facts:

- $A^*$  unique
- $A^*$  exists
- $A^*$  linear
- $A$  bounded  $\Rightarrow A^*$  bounded
- $A$  compact  $\Rightarrow A^*$  compact

- What is the adjoint operator in finite dimensions? (in matrix representation)

- What do you expect to happen with integral operators?
- Adjoint of the single-layer?
- Adjoint of the double-layer?

$$Af(x) = \int k(x,y) f(y) dy$$

$$(Af, g) = \int \int k(x,y) f(y) dy g(x) dx$$

$$= \int \int k(x,y) g(x) dx f(y) dy$$

$$k^*(x,y) = k(y,x)$$

$$= \int \int k(x,y) g(x) dx f(y) dy$$

$$= \int \int k^*(y,x) g(x) dx f(y) dy$$

$$\underbrace{\int \int k^*(y,x) g(x) dx}_{A^*g} = \int_{\Gamma} \log(|x-y|) \rho(y) dy$$

$$= (f, A^*g) \quad \left[ \begin{array}{l} k \rightarrow D\phi(x) = \int \partial_n \log(|x-y|) \rho(y) dy \\ k' \rightarrow S\psi(x) = \int \partial_n \log(|x-y|) \rho(y) dy \end{array} \right]$$

$$D^*\psi(x) = \int \partial_n \log(|x-y|) \rho(y) dy$$

## Fredholm Alternative

**Theorem 23. (Fredholm Alternative [Kress LIE Thm. 4.14])**

$A: X \rightarrow X$  compact. **Then either:**

- $I - A$  and  $I - A^*$  are bijective

or:

- $\dim N(I - A) = \dim N(I - A^*)$

- $(I - A)(X) = N(I - A^*)^\perp$

- $(I - A^*)(X) = N(I - A)^\perp$

$$(I - A)^* = I - A^*$$

$$((I - A)f, y) = (If, y) - (Af, y)$$

$$N(I - A^*) = \text{span}\{1\}$$

- Seen these statements before?
- Translate to language of integral equation solvability:
- What about symmetric kernels ( $K(x, y) = K(y, x)$ )?
- Where to get uniqueness?

$$f \in \text{span}\{1\}^\perp$$

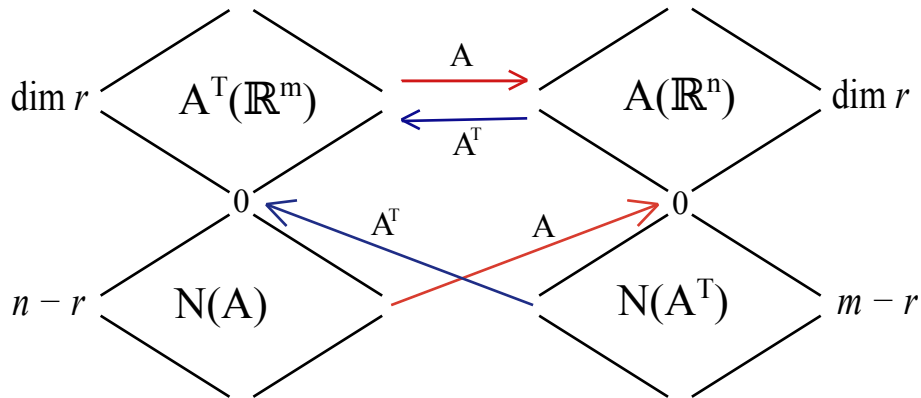
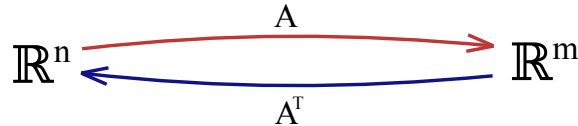
$\Leftrightarrow$

$$(f, 1) = 0 \Leftrightarrow \int f = 0$$

$$(I - A)p = f$$

# Fundamental Theorem of Linear Algebra

$$A \in \mathbb{R}^{m \times n}$$



$$A(x) = N(A^T)^\perp$$

## 6.5 A Tiny Bit of Spectral Theory

## Spectral Theory: Terminology

$A: X \rightarrow X$  bounded,  $\lambda$  is a \_\_\_\_\_ value:

**Definition 24. (Eigenvalue)** *There exists an element  $\phi \in X$ ,  $\phi \neq 0$  with  $A\phi = \lambda\phi$ .*

**Definition 25. (Regular value)** *The "resolvent"  $(\lambda I - A)^{-1}$  exists and is bounded.*

$\rightarrow A\psi = \lambda\psi \quad (\lambda I - A)\psi = 0$

- Can a value be regular and "eigen" at the same time?
- What's special about  $\infty$ -dim here?

**Definition 26. (Resolvent set)**  $\rho(A) := \{\lambda \text{ is regular}\}$

**Definition 27. (Spectrum)**  $\sigma(A) := \mathbb{F}^1 \setminus \rho(A)$

## Spectral Theory of Compact Operators

**Theorem 28.**  $A: X \rightarrow X$  compact linear operator,  $X$   $\infty$ -dim.

**Then:**

- $0 \in \sigma(A)$  (show!)
- $\sigma(A) \setminus \{0\}$  consists *only* of eigenvalues
- $\sigma(A) \setminus \{0\}$  is at most countable
- $\sigma(A)$  has no accumulation point except for 0

- Show first part.  $\rightarrow$  suppose  $0 \notin \sigma(A) \Rightarrow (-A)^{-1} \Rightarrow \underbrace{AA^{-1}}_{\text{compact}} = I$
- Show second part.
- Rephrase last two: how many eigenvalues with  $|\cdot| \geq R$ ?
- **Recap:** What do compact operators do to high-frequency data?
- Don't confuse  $I - A$  with  $A$  itself!

## 7 Singular Integrals and Potential Theory

## Recap: Layer potentials

$$(S\sigma)(x) := \int_{\Gamma} G(x-y) \sigma(y) dS_y$$

$$(S'\sigma)(x) := PV \hat{n} \cdot \nabla_x \int_{\Gamma} G(x-y) \sigma(y) dS_y$$

$$(D\sigma)(x) := PV \int_{\Gamma} \hat{n} \cdot \nabla_y G(x-y) \sigma(y) dS_y$$

$$(D'\sigma)(x) := f.p. \hat{n} \cdot \nabla_x \int_{\Gamma} \hat{n} \cdot \nabla_y G(x-y) \sigma(y) dS_y$$

**Definition 29. (Harmonic function)**  $\Delta u = 0$

- Where are layer potentials harmonic?

## On the double layer again

Is the double layer *actually* weakly singular?

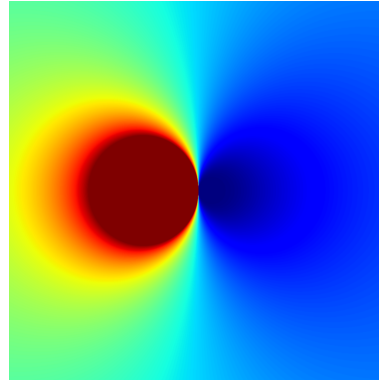
Recap:

### Definition 30. (Weakly singular kernel)

- $K$  defined, continuous everywhere except at  $x = y$
- There exist  $C > 0$ ,  $\alpha \in (0, n - 1]$  such that

$$|K(x, y)| \leq C |x - y|^{\alpha - n + 1} \quad (x, y \in \partial \Omega, x \neq y)$$

$$\frac{\partial}{\partial x} \log(|0 - x|) = \frac{x}{x^2 + y^2}$$



- Singularity with approach on  $y = 0$ ?
- Singularity with approach on  $x = 0$ ?

So life is simultaneously worse and better than discussed.

How about 3D? ( $-x/|x|^3$ )

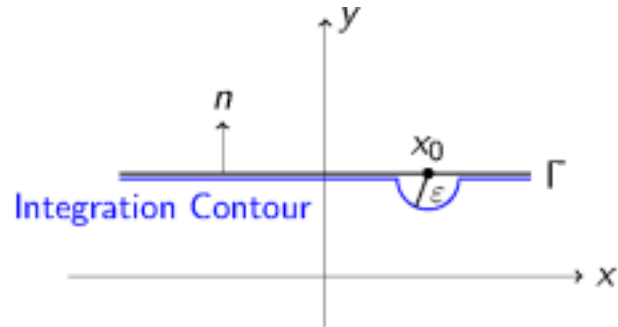
Would like an analytical tool that requires 'less' fanciness.

## Cauchy Principal Value

- But I don't **want** to integrate across a singularity!
- 

$$\int_{-1}^1 \frac{1}{x} dx?$$

## Principal Value in $n$ dimensions



- Again: Symmetry matters!
- What about even worse singularities?

## Recap: Layer potentials

$$(S\sigma)(x) := \int_{\Gamma} G(x-y) \sigma(y) \, ds_y$$

$$(S'\sigma)(x) := \text{PV} \hat{n} \cdot \nabla_x \int_{\Gamma} G(x-y) \sigma(y) \, ds_y$$

$$(D\sigma)(x) := \text{PV} \int_{\Gamma} \hat{n} \cdot \nabla_y G(x-y) \sigma(y) \, ds_y$$

$$(D'\sigma)(x) := \text{f.p.} \hat{n} \cdot \nabla_x \int_{\Gamma} \hat{n} \cdot \nabla_y G(x-y) \sigma(y) \, ds_y$$

**Important for us:** Recover 'average' of interior and exterior limit without having to refer to off-surface values.

## Green's Theorem

**Theorem 31. (Green's Theorem [Kress LIE Thm 6.3])**

$$\int_{\Omega} u \Delta v + \nabla u \cdot \nabla v = \int_{\partial\Omega} u (\hat{n} \cdot \nabla v) \, dS$$

$$\int_{\Omega} u \Delta v - v \Delta u = \int_{\partial\Omega} u (\hat{n} \cdot \nabla v) - v (\hat{n} \cdot \nabla u) \, dS$$

- If  $\Delta v = 0$ , then

$$\int_{\partial\Omega} \hat{n} \cdot \nabla v = ?$$

- What if  $\Delta v = 0$  and  $u = G(|y - x|)$  in Green's second identity?

## Green's Formula

**Theorem 32. (Green's Formula [Kress LIE Thm 6.5])** *If  $\Delta u = 0$ , then*

$$(S(\hat{n} \cdot \nabla u) - Du)(x) = \begin{cases} u(x) & x \in D \\ \frac{u(x)}{2} & x \in \partial D \\ 0 & x \notin D \end{cases}$$

- Suppose I know 'Cauchy data' ( $u|_{\partial D}, \hat{n} \cdot \nabla u|_{\partial D}$ ) of  $u$ . What can I do?
- What if  $D$  is an exterior domain?

## Things harmonic functions (don't) do

**Theorem 33. (Mean Value Theorem [Kress LIE Thm 6.7])** If  $\Delta u = 0$ ,

$$u(x) = \int_{H(x,r)} u(y) \, d\bar{y} = \int_{\partial H(x,r)} u(y) \, d\bar{y}$$

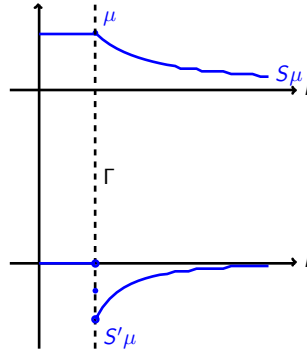
- Define  $\int \bar{\cdot}$  ?
- Trace back to Green's Formula (say, in 2D):

**Theorem 34. (Maximum Principle [Kress LIE 6.9])** If  $\Delta u = 0$  on compact set  $\bar{D}$ :

*$u$  attains its maximum on the boundary.*

- Suppose it were to attain its maximum somewhere inside an open set...
- What do our **constructed** harmonic functions (i.e. layer potentials) do there?

## Jump relations



Let  $[X] = X_+ - X_-$ . (Normal points towards "+"="exterior".)

[Kress LIE Thm. 6.14, 6.17, 6.18]

$$\lim_{x \rightarrow x_0 \pm} (S' \sigma) = \left( S' \mp \frac{1}{2} I \right) (\sigma)(x_0) \quad \Rightarrow \quad [S' \sigma] = -\sigma$$

$$\lim_{x \rightarrow x_0 \pm} (D \sigma) = \left( D \pm \frac{1}{2} I \right) (\sigma)(x_0) \quad \Rightarrow \quad [D \sigma] = \sigma$$

$$[D'\sigma]=0$$

- Truth in advertising: Assumptions on  $\Gamma$ ?

## Green's Formula at Infinity (skipped)

$\Omega \subseteq \mathbb{R}^n$  bounded,  $C^1$ , connected boundary,  $\Delta u = 0$ ,  $u$  bounded

$$(\mathcal{S}_{\partial\Omega}(\hat{n} \cdot \nabla u) - D_{\partial\Omega} u)(x) + (\mathcal{S}_{\partial B_r}(\hat{n} \cdot \nabla u) - D_{\partial B_r} u)(x) = u(x)$$

for  $x$  between  $\partial\Omega$  and  $B_r$ .

Now  $r \rightarrow \infty$ .

Behavior of individual terms?

**Theorem 35. (Green's Formula in the exterior [Kress LIE Thm 6.10])**

$$(\mathcal{S}_{\partial\Omega}(\hat{n} \cdot \nabla u) - D_{\partial\Omega} u)(x) + \text{PV } u_\infty = u(x)$$

for some constant  $u_\infty$ . *Only* for  $n=2$ ,

$$u_\infty = \frac{1}{2\pi r} \int_{|y|=r} u(y) dS_y.$$

**Theorem 36. (Green's Formula in the exterior [KRESS LIE Thm 6.10])**

$$(S_{\partial\Omega}(\hat{n} \cdot \nabla u) - D_{\partial\Omega} u)(x) + u_{\infty} = u(x)$$

- Realize the power of this statement:
- Behavior of the fundamental solution as  $r \rightarrow \infty$ ?
- How about its derivatives?

## 8 Boundary Value Problems

### 8.1 Laplace

## Boundary Value Problems: Overview

	Dirichlet	Neumann
<b>Int.</b>	$\lim_{x \rightarrow \partial\Omega^-} u(x) = g$ + unique	$\lim_{x \rightarrow \partial\Omega^-} \hat{n} \cdot \nabla u(x) = g$ ⦿ may differ by constant
<b>Ext.</b>	$\lim_{x \rightarrow \partial\Omega^+} u(x) = g$ $u(x) = \begin{cases} o(1) & \frac{2}{D} \\ o(1) & \frac{3}{D} \end{cases}$ as $ x  \rightarrow \infty$ + unique	$\lim_{x \rightarrow \partial\Omega^+} \hat{n} \cdot \nabla u(x) = g$ $u(x) = o(1)$ as $ x  \rightarrow \infty$ + unique

with  $g \in C(\partial\Omega)$ .

- What does  $f(x) = O(1)$  mean? (and  $f(x) = o(1)$ ?)
- Dirichlet uniqueness: why?
- Neumann uniqueness: why?
- Truth in advertising: Missing assumptions on  $\Omega$ ?
- What's a DtN map?

**Next mission:** Find IE representations for each.

## Uniqueness of Integral Equation Solutions

### Theorem 37. (Nullspaces [Kress LIE Thm 6.20])

- $\mathcal{N}(I/2 - D) = \mathcal{N}(I/2 - S') = \{0\}$
- $\mathcal{N}(I/2 + D) = \text{span}\{1\}$ ,  $\mathcal{N}(I/2 + S') = \text{span}\{\psi\}$ ,  
where  $\int \psi \neq 0$ .

- Show  $\mathcal{N}(I/2 - D) = \{0\}$ .
  - Show  $\mathcal{N}(I/2 - S') = \{0\}$ .
  - Show  $\mathcal{N}(I + D) = \text{span}\{1\}$ .
  - What extra conditions on the RHS do we obtain?
- “Clean” Existence for 3 out of 4.

## Patching up Exterior Dirichlet (skipped)

Problem:  $\mathcal{N}(I/2 + S') = \{\psi\}$ ...but we do not know  $\psi$ .

Use a different kernel:

$$\hat{n} \cdot \nabla_y G(x, y) \quad \rightarrow \quad \hat{n} \cdot \nabla_y G(x, y) + \frac{1}{|x|^{n-2}}$$

Note: Singularity only at origin! (assumed  $\in \Omega$ )

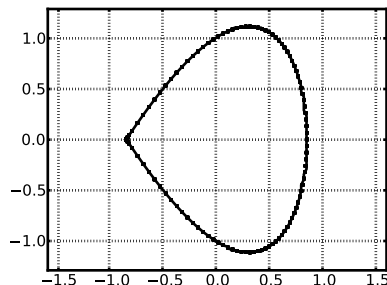
- 2D behavior? 3D behavior?
- Still a solution of the PDE?
- Compact?
- Jump condition? Exterior limit? Deduce  $u = 0$  on exterior.
- $|x|^{n-2} u(x) = ?$  on exterior
- Thus  $\int \phi = 0$ . Contribution of the second term?
- $\phi/2 + D\phi = 0$ , i.e.  $\phi \in \mathcal{N}(I/2 + D) = ?$

- Existence/uniqueness?

→ Existence for 4 out of 4.

- Remaining key shortcoming of IE theory for BVPs?

## Domains with Corners



What's the problem? (Hint: Jump condition for constant density)

At corner  $x_0$ : (2D)

$$\lim_{x \rightarrow x_0 \pm} = \int_{\partial\Omega} \hat{n} \cdot \nabla_y G(x, y) \phi(y) d^2s_y \pm \frac{1}{2} \frac{\langle \text{opening angle on } \pm \text{ side} \rangle}{\pi} \phi$$

→ non-continuous behavior of potential on  $\Gamma$  at  $x_0$

What space have we been living in?

Fixes:

- $L^2$  + Bounded (Neumann) + Compact (Fredholm)

- Use  $L^2$  theory  
(point behavior “invisible”)

Numerically: Needs consideration, but ultimately easy to fix.

## 8.2 Helmholtz

8.2.1 Helmholtz

8.2.2 Helmholtz

8.2.3 Helmholtz

8.2.4 Helmholtz

8.2.5 Helmholtz

8.2.6 Helmholtz

8.2.7 Helmholtz

8.2.8 Helmholtz

8.2.9 Helmholtz

8.2.10 Helmholtz

8.2.11 Helmholtz

8.2.12 Helmholtz

8.2.13 Helmholtz

8.2.14 Helmholtz

8.2.15 Helmholtz

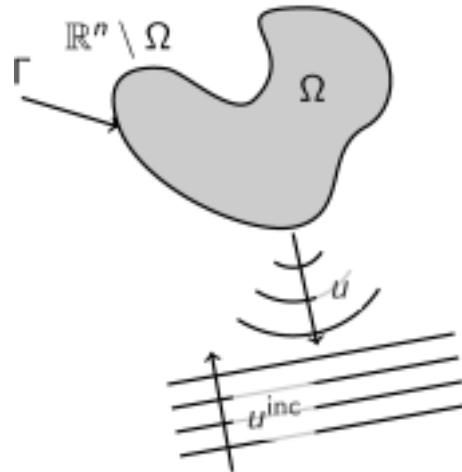
8.2.16 Helmholtz

8.2.17 Helmholtz

8.2.18 Helmholtz

8.2.19 Helmholtz

## The prototypical Helmholtz BVP: A Scattering Problem



Ansatz:

$$u^{\text{tot}} = u + u^{\text{inc}}$$

Solve for scattered field  $u$ .

## Helmholtz: Some Physics

Physical quantities:

- Velocity potential:  $U(x, t) = u(x) e^{-i\omega t}$   
(fix phase by e.g. taking real part)
- Velocity:  $v = (1/\rho_0) \nabla U$
- Pressure:  $p = -\partial_t U = i\omega u e^{-i\omega t}$ 
  - Equation of state:  $p = f(\rho)$
- What's  $\rho_0$ ?
- What happens to a pressure BC as  $\omega \rightarrow 0$ ?

## Helmholtz: Boundary Conditions

- **Sound-soft:** Pressure remains constant
  - Scatterer “gives”
  - $u = f \rightarrow$  Dirichlet
- **Sound-hard:** Pressure same on both sides of interface
  - Scatterer “does not give”
  - $\hat{n} \cdot \nabla u = 0 \rightarrow$  Neumann
- **Impedance:** Some pressure translates into motion
  - Scatterer “resists”
  - $\hat{n} \cdot \nabla u + ik\lambda u = 0 \rightarrow$  Robin ( $\lambda > 0$ )
- **Sommerfeld** radiation condition: allow only outgoing waves

$$r^{\frac{n-1}{2}} \left( \frac{\partial}{\partial r} - ik \right) u(\mathbf{x}) \rightarrow 0 \quad (r \rightarrow \infty)$$

Many interesting BCs → many IEs! :)

- Transmission between media: What's continuous?

## Unchanged from Laplace

**Theorem 38. (Green's Formula [Colton/Kress IAST Thm 2.1])** If  $\Delta u + k^2 u = 0$ , then

$$(S(\hat{n} \cdot \nabla u) - Du)(x) = \begin{cases} u(x) & x \in D \\ \frac{u(x)}{2} & x \in \partial D \\ 0 & x \notin D \end{cases}$$

$$\begin{aligned} \lim_{x \rightarrow x_0 \pm} (S' u) &= \left( S' \mp \frac{1}{2} I \right) (u)(x_0) & \Rightarrow & [S u] = 0 \\ & & & [S' u] = -u \\ \lim_{x \rightarrow x_0 \pm} (D u) &= \left( D \pm \frac{1}{2} I \right) (u)(x_0) & \Rightarrow & [D u] = u \\ & & & = 0 \end{aligned}$$

- Why is singular behavior (esp. jump conditions) unchanged?
- Why does Green's formula survive?

Remember Green's theorem:

$$\int_{\Omega} u \Delta v - v \Delta u = \int_{\partial\Omega} u (\hat{n} \cdot \nabla v) - v (\hat{n} \cdot \nabla u) \, d\mathbf{s}$$

## Resonances

–  $\Delta$  on a bounded (interior) domain with homogeneous Dirichlet/Neumann BCs has countably many real, positive eigenvalues.

- What does that have to do with Helmholtz?
- Why could it cause grief?

## Helmholtz: Boundary Value Problems

Find  $u \in C(\bar{D})$  with  $\Delta u + k^2 u = 0$  such that

	Dirichlet	Neumann
Int.	$\lim_{x \rightarrow \partial D^-} u(x) = g$ ⦿ unique (−resonances)	$\lim_{x \rightarrow \partial D^-} \hat{n} \cdot \nabla u(x) = g$ ⦿ unique (−resonances)
Ext.	$\lim_{x \rightarrow \partial D^+} u(x) = g$ Sommerfeld ⦿ unique	$\lim_{x \rightarrow \partial D^+} \hat{n} \cdot \nabla u(x) = g$ Sommerfeld ⦿ unique

with  $g \in C(\partial D)$ .

- Find layer potential representations for each.

## Patching up resonances

**Issue:** Ext. IE inherits non-uniqueness from 'adjoint' int. BVP

**Fix:** Tweak representation [Brakhage/Werner '65, ...]  
(also called the 'CFIE'—'combined field integral equation')

$$u = D\phi - i\alpha S\phi$$

( $\alpha$ : tuning knob  $\rightarrow 1$  is fine,  $\sim k$  better for large  $k$ )

- How does this help? (skipped)
- Uniqueness for remaining IEs similar. (skipped)

## 8.3 Calderón identities

- Show that  $D'$  is self adjoint.
- Show that  $(S\varphi, D'\psi) = ((S' + I/2)\varphi, (D - I/2)\psi)$ .
- $(\varphi, S D'\psi)$ ?

## Calderón Identities: Summary

- $S D' = D^2 - I/4$
- $D' S = S'^2 - I/4$

Also valid for Laplace (jump relation same after all!)

- Why do we care?

(to end lec 18)

## 9 Back from Infinity: Discretization

# 10 Computing Integrals: Approaches to Quadrature

## 11 Going General: More PDEs